# Numerical Fourier analysis of quasi-periodic functions 

G. Gómez, ${ }^{1} \quad$ J.M. Mondelo ${ }^{2} \quad$ C. Simó ${ }^{1}$<br>${ }^{1}$ Departament de Matemàtica Aplicada i Anàlisi, Universitat de Barcelona<br>${ }^{2}$ Departament de Matemàtiques, Universitat Autònoma de Barcelona<br>\section*{WSIMS08}<br>IMUB dec 1-5, 2008

## Outline

Introduction

The method

Error estimation

Accuracy test

Study of the stability region around $L_{5}$

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## Introduction

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## Setting

We are given an analytic, quasi-periodic function

$$
f(t)=\sum_{\boldsymbol{k} \in \mathbb{Z}^{m}} a_{\boldsymbol{k}} e^{i 2 \pi\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle t}
$$

satisfying the Cauchy estimates

$$
\left|a_{k}\right| \leq C e^{-\delta|\boldsymbol{k}|} \quad\left(\exists C>0, \delta>0, \quad|\boldsymbol{k}|=\left|\left(k_{1}, \ldots, k_{m}\right)\right|=\left|k_{1}\right|+\cdots+\left|k_{m}\right|\right)
$$

and with a vector of basic frequencies $\boldsymbol{\omega}=\left(\omega_{1}, \ldots, \omega_{m}\right)$ satisfying a Diophantine condition

$$
|\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle|>\frac{D}{|k|^{\tau}}, \quad(\exists D, \tau>0)
$$

We want to numerically compute the frequencies $\{\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle\}_{|\boldsymbol{k}|=0}^{\text {maxor }}$ and amplitudes $a_{k}$ from the values of $f$.

## Fourier Transform

The Fourier Transform will be denoted as

$$
f(t) \quad \xrightarrow{\mathcal{F}} \quad \mathcal{F}(f(t))(\omega)=\hat{f}(\omega)=\int_{-\infty}^{\infty} f(t) e^{-i 2 \pi \omega t} d t
$$

If $f(t)$ is quasi-periodic, its Fourier transform is a discrete set of impulses based at the frequencies:

$$
\begin{aligned}
& f(t)=\sum_{\boldsymbol{k} \in \mathbb{Z}^{m}} a_{\boldsymbol{k}} e^{i 2 \pi\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle t} \quad \stackrel{\mathcal{F}}{\longrightarrow} \hat{f}(\omega)=\sum_{\boldsymbol{k} \in \mathbb{Z}^{m}} a_{\boldsymbol{k}} \delta_{\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle}(\omega)
\end{aligned}
$$

Example: $f(t)=\cos (2 \pi 0.1 t)+0.5 \cos (2 \pi 0.2 t)+0.4 \cos (2 \pi 0.35 t)$

## Time truncation $\longrightarrow$ WFT

Graphical development (E.O. Brigham, 1988)
Time truncation gives rise to the phenomenon known as leakage.
Example: $T=40, f(t)=\cos (2 \pi 0.1 t)+0.5 \cos (2 \pi 0.2 t)+0.4 \cos (2 \pi 0.35 t)$.



$\downarrow \mathcal{F}$


$\downarrow \mathcal{F}$


The maxima of the WFT (bottom right) are displaced from the true frequencies.

## Time truncation $\longrightarrow$ WFT

Explicit formulae

- Windowed Fourier Transform:

$$
\begin{aligned}
\phi_{f, T}(\omega) & :=\frac{1}{T} \mathcal{F}\left(\chi_{[0, T]} f(t)\right)(\omega) \\
& =\frac{1}{T} \int_{0}^{T} \chi_{[0, T]}(t) f(t) e^{-i 2 \pi \omega t} d t
\end{aligned}
$$

- Leakage of a complex exponential term.

$$
\begin{aligned}
\left|\phi_{e^{i 2 \pi \nu t}, T}(\omega)\right| & =\left|\frac{e^{i 2 \pi(\nu-\omega) T}-1}{i 2 \pi(\nu-\omega) T}\right| \\
& =\left|\frac{\sin \pi(\nu-\omega) T}{\pi(\nu-\omega) T}\right| \\
& =|\operatorname{sinc}((\nu-\omega) T)|
\end{aligned}
$$

## Reducing leackage

There are two strategies:

- Increase the window length.

$$
\left|\phi_{e^{i 2 \pi \nu t}, T}(\omega)\right|=|\operatorname{sinc}((\nu-\omega) T)|=\left|\frac{\sin \pi(\nu-\omega) T}{\pi(\nu-\omega) T}\right|
$$




## Reducing leackage

There are two strategies:

- Use a smoother window. We use Hanning's:

$$
H_{T}^{n_{h}}(t)=q_{n_{h}}\left(1-\cos \frac{2 \pi t}{T}\right)^{n_{h}}
$$

being $q_{n_{h}}=n_{h}!/\left(\left(2 n_{h}-1\right)!!\right)$.
The corresponding WFT is denoted by

$$
\phi_{f, T}^{n_{h}}(\omega):=\mathcal{F}\left(H^{n_{h}} f\right)(\omega)=\frac{1}{T} \int_{0}^{T} H_{T}^{n_{h}}(t) f(t) e^{-i 2 \pi \omega t} d t
$$

## Reducing leackage

There are two strategies:

- Use a smoother window.

$$
\phi_{e^{i 2 \pi \nu t}, T}(\omega)=\frac{e^{i 2 \pi(\nu-\omega) T}-1}{i 2 \pi(\nu-\omega) T}=O\left(\frac{1}{(\nu-\omega) T}\right)
$$

vS

$$
\phi_{e^{12 \pi \nu t}, T}^{n_{h}}(\omega)=\frac{(-1)^{n_{h}}\left(n_{h}!\right)^{2}\left(e^{i 2 \pi(\nu-\omega) T}-1\right)}{i 2 \pi \prod_{j=-n_{h}}^{n_{h}}((\nu-\omega) T+j)}=O\left(\frac{1}{((\nu-\omega) T)^{1+2 n_{h}}}\right)
$$




## Discretization $\longrightarrow$ DFT

Graphical development (E.O. Brigham, 1988)

$$
\begin{aligned}
& T=40, N=32, f(t)=\cos (2 \pi 0.1 t)+0.5 \cos (2 \pi 0.2 t)+0.4 \cos (2 \pi 0.35 t) \\
& \times \\
& \text { * }
\end{aligned}
$$


impulse spacing $=$ sampling rate $=\mathrm{T} / \mathrm{N}=1.25$



impulse spacing $=$ DFT period $=\mathrm{N} / \mathrm{T}=0.8$

$\downarrow$


## Sampling $\longrightarrow$ DFT

Explicit formulae

- DFT of $\left\{f\left(j \frac{T}{N}\right)\right\}_{j=0}^{N-1}$ defined as $\left\{F_{f, T, N}(k)\right\}_{k=0}^{N-1}$, being

$$
\begin{aligned}
F_{f, T, N}(k) & :=\frac{1}{N} \mathcal{F}\left(\sum_{j \in \mathbb{Z}} \chi_{[0, T]}\left(j \frac{T}{N}\right) f\left(j \frac{T}{N}\right) \delta_{j^{T}}\right. \\
& =\frac{k}{N} \sum_{j=0}^{N} f\left(j \frac{T}{N}\right) e^{-i 2 \pi k j / N}
\end{aligned}
$$

- With Hanning's window:

$$
F_{f, T, N}^{n_{h}}(k)=\frac{1}{N} \sum_{j=0}^{N-1} H_{T}^{n_{h}}\left(j \frac{T}{N}\right) f\left(j \frac{T}{N}\right) e^{-i 2 \pi k j / N}
$$

## Sampling $\longrightarrow$ DFT

Explicit formulae

- Relation with the WFT:

$$
F_{f, T, N}(k)=\phi_{f, T, N}\left(\frac{k}{T}\right)+\underbrace{\sum_{l \in \mathbb{Z} \backslash\{0\}}\left(\phi_{f, T, N}\left(\frac{k+l N}{T}\right)+\phi_{f, T, N}\left(\frac{k-l N}{T}\right)\right)}_{\text {error term }}
$$

- The fundamental domain of the DFT for real signals is $[0, T /(2 N)]$. $T /(2 N)$ is Nyquist's critical frequency.



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- The fundamental domain of the DFT for real signals is $[0, T /(2 N)]$. $T /(2 N)$ is Nyquist's critical frequency.
- The error term above can produce aliasing: if a frequency of the signal is outside the fundamental domain of the DFT, we will detect an alias of it.
- Aliasing is avoided increasing $N$.


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## Algorithm

Parameters: $T$ (time length), $N$ (number of samples), $n_{h}$ (Hanning index) $b_{\text {min }}$ minimum threshold, several tolerances.

1. Set an starting threshold for collecting peaks of the modulus of the DFT of $f(t)$.
2. Find initial approximations of the frequencies, starting from the peaks of the DFT greater than the thresold.
3. Find the amplitudes of the frequencies found in the previous step, by solving $\operatorname{DFT}\left(Q_{f}\right)=\operatorname{DFT}(f)$
4. Simultaneously refine ALL the frequencies and amplitudes of the current quasi-periodic approximation of $f$, by solving $\operatorname{DFT}\left(Q_{f}\right)=\operatorname{DFT}(f)$.
5. Perform a DFT of the input signal minus the current quasi-periodic approximation obtained in step 4, decrease the thresold and go back to step 2.

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## An illustration of the algorithm

For $f(t)=\cos (2 \pi 0.13 t)-\frac{1}{2} \sin (2 \pi 0.27 t)+\sin (2 \pi 0.37 t)$, $T=N=512, n_{h}=0$.

1. Starting thresold: 0.8 modulus of the DFT of the input data:

$\Rightarrow$ peaks $j=61, j=189$.

## An illustration of the algorithm

For $f(t)=\cos (2 \pi 0.13 t)-\frac{1}{2} \sin (2 \pi 0.27 t)+\sin (2 \pi 0.37 t)$, $T=N=512, n_{h}=0$.
2. Approximation of frequencies:

$$
\begin{aligned}
\text { peak } 67 & \Rightarrow \text { frequency } 0.130859375 \\
\text { peak } 189 & \Rightarrow \text { frequency } 0.369140625
\end{aligned}
$$

3. Computation of amplitudes from known frequencies:

| Frequency | Cosine amplitude | Sine amplitude |
| :--- | :--- | :--- |
| 0.369140625 | 0.702312716711 | 0.136800713691 |
| 0.130859375 | 0.137731069235 | 0.699288924190 |

modulus of the DFT of the residual


## An illustration of the algorithm

For $f(t)=\cos (2 \pi 0.13 t)-\frac{1}{2} \sin (2 \pi 0.27 t)+\sin (2 \pi 0.37 t)$, $T=N=512, n_{h}=0$.
4. Iterative refinement:

| Frequency | Cosine amplitude | Sine amplitude |
| :--- | :--- | ---: |
| 0.369995932915 | 0.005462459021 | 1.000450861577 |
| 0.129998625183 | 0.999908805689 | -0.002241420351 |

5. modulus of the DFT of input signal minus step 4:



New threshold: 0.2

## An illustration of the algorithm

For $f(t)=\cos (2 \pi 0.13 t)-\frac{1}{2} \sin (2 \pi 0.27 t)+\sin (2 \pi 0.37 t)$, $T=N=512, n_{h}=0$.
5. modulus of the DFT of input signal minus step 4:


New threshold: 0.2
2. Approximation of frequencies:

$$
\text { peak } 138 \Rightarrow \text { frequency } 0.26953125
$$

3. Amplitudes from known frequencies:

| Frequency | Cosine amplitude | Sine amplitude |
| :--- | ---: | :---: |
| 0.369995932915 | 0.005462459021 | 1.000450861577 |
| 0.129998625183 | 0.999908805689 | -0.002241420352 |
| 0.269531250000 | -0.309714556917 | -0.330986794067 |

## An illustration of the algorithm

For $f(t)=\cos (2 \pi 0.13 t)-\frac{1}{2} \sin (2 \pi 0.27 t)+\sin (2 \pi 0.37 t)$, $T=N=512, n_{h}=0$.
4. Iterative refinement:

| Frequency | Cosine amplitude | Sine amplitude |
| :--- | ---: | ---: |
| 0.3700000000000000 | 0.0000000000000009 | 1.0000000000000022 |
| 0.1300000000000000 | 0.9999999999999997 | 0.0000000000000010 |
| 0.2700000000000000 | -0.0000000000000028 | -0.4999999999999995 |

modulus of the DFT of the residual:


## Computing amplitudes from known frequencies

We ask $\operatorname{DFT}\left(Q_{f}\right)=\operatorname{DFT}(f)$, being

$$
Q_{f}(t)=A_{0}^{c}+\sum_{l=1}^{N_{f}}\left(A_{l}^{c} \cos \left(2 \pi \frac{\nu_{l}}{T} t\right)+A_{l}^{s} \sin \left(2 \pi \frac{\nu_{l}}{T} t\right)\right.
$$

Since we work with real signals, we use the sine and cosine transforms:

$$
\begin{array}{ll}
c_{f, T, N}^{n_{h}}(k)=\frac{2}{N} \sum_{j=0}^{N-1} f\left(j \frac{T}{N}\right) H_{N}^{n_{h}}(j) \cos \left(2 \pi \frac{k}{N} j\right), \quad k=0, \ldots, \frac{N}{2}, \\
s_{f, T, N}^{n_{h}}(k)=\frac{2}{N} \sum_{j=0}^{N-1} f\left(j \frac{T}{N}\right) H_{N}^{n_{h}}(j) \sin \left(2 \pi \frac{k}{N} j\right), \quad k=1, \ldots, \frac{N}{2}-1 .
\end{array}
$$

They are realted to the DFT in complex form by

$$
F_{f, T, N}^{n_{h}}(k)=\frac{1}{2}\left(c_{f, T, N}^{n_{h}}(k)-i s_{f, T, N}^{n_{h}}(k)\right), \quad k=0, \ldots, N / 2 .
$$

## Computing amplitudes from known frequencies

The system of equations to be solved is linear and $\left(1+2 N_{f}\right) \times\left(1+2 N_{f}\right)$ :

$$
\begin{aligned}
& A_{0}^{c} c_{1, T, N}^{n_{h}}(0)+\sum_{\substack{ \\
N_{f}}}^{N_{N_{f}}}\left(A_{l}^{c} \bar{c}_{\nu_{l}, N}^{n_{h}}(0)+A_{l}^{s} c_{\nu_{l}, N}^{n_{h}}(0)\right)=c_{f, T, N}^{n_{h}}(0) \\
& \begin{aligned}
A_{0}^{c} c_{1, T, N}^{n_{h}}(j)+\sum_{l=1}^{N_{f}}\left(A_{l}^{c} \bar{c}_{\nu_{l}, N}^{n_{h}}(j)+A_{l}^{s} \widetilde{c}_{\nu_{l}, N}^{n_{h}}(j)\right) & =c_{f, T, N}^{n_{h}}(j) \\
\sum_{l=1}^{N_{f}}\left(A_{l}^{c} \bar{s}_{\nu_{l}, T}^{n_{h}}(j)+A_{l}^{c} s_{\nu_{l}, T}^{n_{h}}(j)\right) & =s_{f, T, N}^{n_{h}}(j)
\end{aligned}
\end{aligned}
$$

where $j=\left[\nu_{l}+0.5\right], l=1 \div N_{f}$ (collocation harmonics), and

$$
\begin{aligned}
c_{1}^{n_{h}}(j) & =c_{1, T, N}^{n_{h}}(j), \\
\bar{c}_{\nu_{l}, N}^{n_{h}}(j) & =c_{\cos \left(\frac{2 \pi \nu_{l}}{T}\right), T, N}^{n_{h}}(j), \quad \bar{s}_{\nu_{l}, N}^{n_{h}}(j)=s_{\cos \left(\frac{2 \pi \nu_{l}}{T}\right), T, N}^{n_{h}}(j), \\
\widetilde{c}_{\nu_{l}, N}^{n_{h}}(j) & =c_{\sin \left(\frac{2 \pi \nu_{l}}{T}\right), T, N}^{n_{h}}(j), \quad \tilde{s}_{\nu_{l}, N}^{n_{h}}(j)=s_{\sin \left(\frac{2 \pi \nu_{l}}{T}\right), T, N}^{n_{h}}(j) .
\end{aligned}
$$

## Simultaneous improvement of frequencies and amplitudes

We solve by Newton's method the following $\left(1+3 N_{f}\right) \times\left(1+3 N_{f}\right)$ non-linear system:

$$
\begin{aligned}
& A_{0}^{c} c_{1, T, N}^{n_{h}}(0)+\sum_{\substack{l=1 \\
N_{f}}}^{N_{f}}\left(A_{l}^{c} \bar{c}_{\nu_{l}, N}^{n_{h}}(0)+A_{l}^{s} c_{\nu_{l}, N}(0)\right)=c_{f, T, N}^{n_{h}}(0) \\
& \begin{array}{r}
A_{0}^{c} c_{1, T, N}^{n_{h}}\left(j_{i}\right)+\sum_{\substack{l_{1} \\
N_{f}}}\left(A_{l}^{c} c_{\nu_{l, N}}^{n_{h}}\left(j_{i}\right)+A_{l}^{s} c_{\nu_{l, N}}^{n_{h}}\left(j_{i}\right)\right)
\end{array}=c_{f, T, N}^{n_{h}}\left(j_{i}\right) . \\
& A_{0}^{c} c s_{1, T, N}^{n_{n}}\left(j_{i}^{+}\right)+\sum_{l=1}^{N_{f}}\left(A_{l}^{c} \widetilde{c s}_{\nu_{l}, N}^{n_{h}}\left(j_{i}^{+}\right)+A_{l}^{s} \widetilde{c s}_{\nu_{l}, N}^{n_{h}}\left(j_{i}^{+}\right)\right)=c s_{f, T, N}^{n_{h}}\left(j_{i}^{+}\right)
\end{aligned}
$$

being $j_{i}=\left[\nu_{i}+0.5\right], j_{i}^{+}=\left[\nu_{i}\right]+1-\left(j_{i}^{+}-\left[\nu_{i}\right]\right)$.

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## Strategy

Let us denote

- $f_{r_{0}}$ : the truncation of $f$ to the frequencies we want to determine:

$$
f_{r_{0}}(t)=A_{0}^{c}+\sum_{\substack{|k| \leq r_{0}-1 \\\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle>0}}\left(A_{k}^{c} \cos (2 \pi\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle t)+A_{k}^{s} \sin (2 \pi\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle t)\right) .
$$

- $y=\left(A_{0}, \nu_{1}, A_{1}^{c}, A_{1}^{s}, \ldots, \nu_{N_{f}}, A_{N_{f}}^{c}, A_{N_{f}}^{s}\right)$ : the exact frequencies and amplitudes.
- $y+\Delta y$ : the computed frequencies and amplitudes.

The system we solve for iterative improvement of frequencies and amplitudes is

$$
\underbrace{\operatorname{DFT}\left(Q_{f}\right)}_{g(y+\Delta y)}=\underbrace{\operatorname{DFT}\left(f_{r_{0}}\right)}_{b}+\underbrace{\operatorname{DFT}\left(f-f_{r_{0}}\right)}_{\Delta b}
$$

We would get the exact frequencies and amplitudes if $\Delta b=0$.

## Strategy

- System for iterative improvement of frequencies and amplitudes:

$$
\begin{aligned}
& A_{0}^{c} c_{1}^{n_{h}}\left(j_{i}\right)+\sum_{l=1}^{N_{f}}\left(A_{l}^{c} c_{\nu_{l}, N}^{n_{h}}\left(j_{i}\right)+A_{l}^{s} c_{\nu_{l}, N}^{n_{h}}\left(j_{i}\right)\right)=c_{f_{r_{0}}, T, N}^{n_{h}}\left(j_{i}\right)+c_{f-f_{r_{0}, T, N}}^{n_{h}}\left(j_{i}\right) \\
& \sum_{l=1}^{N_{f}}\left(A_{l}^{c \bar{s}_{\nu_{l}, N}^{n_{h}}}\left(j_{i}\right)+A_{l}^{s \sim s_{\nu_{l}} n_{h}}\left(j_{i}\right)\right)=s_{f_{r_{0}}, T, N}^{n_{h}}\left(j_{i}\right)+s_{f-f_{r_{0}}, T, N}^{n_{h}}\left(j_{i}\right) \\
& A_{0}^{c} c s_{1}^{n_{h}}\left(j_{i}^{+}\right)+\sum_{l=1}^{N_{f}}\left(A_{l}^{c} \overline{c s_{\nu_{l}, N}} n_{h}\left(j_{i}^{+}\right)+A_{l}^{s} \widetilde{c} s_{\nu_{l}, N}^{n_{h}}\left(j_{i}^{+}\right)\right)=c s_{f_{r_{0}}, T, N}^{n_{h}}\left(j_{i}^{+}\right)+c s_{f-f_{r_{0}}, T, N}^{n_{h}}\left(j_{i}^{+}\right) \text {. } \\
& \text { where } f-f_{r_{0}}=\sum_{|k| \geq r_{0}} a_{k} e^{i 2 \pi\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle t} \text {. }
\end{aligned}
$$

- The error term $\Delta b$ consists of DFT
- of periodic terms with frequencies not being computed,
- evaluated in harmonics corresponding to frequencies being computed.

Therefore, the error term $\Delta b$ can be considered leakage of the remainder, $f-f_{r_{0}}$.

## Strategy

- The error term $\Delta b$ can be considered leakage of the remainder

$$
\operatorname{DFT}\left(f-f_{r_{0}}\right)=\sum_{|k| \geq r_{0}} a_{k} \operatorname{DFT}\left(e^{i 2 \pi\langle\omega, k\rangle t}\right)
$$

- The effect of the terms of the remainder on the error $\Delta b$ is
- The DFT of terms corresponding to low-order frequencies, $\{\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle\}_{|k| \gtrsim r_{0}}$, evaluated at the harmonics $\left\{j_{i}, j_{i}^{+}\right\}$, will be small if the harmonics $T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle$ are far from $\left\{j_{i}, j_{i}^{+}\right\}$.
This can be achieved by increasing $T$ as long as there is no aliasing.
- The DFT of terms corresponding to high-order frequencies may not be small ( $T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle$ can be made arbitrarily close to a $j_{i}$ for large enough $|\boldsymbol{k}|$ ). However, the corresponding amplitudes will be small due to the Cauchy estimates

$$
\left|a_{k}\right| \leq C e^{-\delta|\boldsymbol{k}|} \quad \forall \boldsymbol{k} \in \mathbb{Z}^{m}
$$

so they will be harmless.

## Bounding

- The system we solve for iterative improvement of frequencies and amplitudes is

$$
\underbrace{\operatorname{DFT}\left(Q_{f}\right)}_{g(y+\Delta y)}=\underbrace{\operatorname{DFT}\left(f_{r_{0}}\right)}_{b}+\underbrace{\operatorname{DFT}\left(f-f_{r_{0}}\right)}_{\Delta b}
$$

We would get the exact frequencies and amplitudes if $\Delta b=0$.

- The error in frequencies and amplitudes is given, at first order, by

$$
\|\Delta y\|_{\infty} \leq\left\|D g(y)^{-1}\right\|_{\infty}\|\Delta b\|_{\infty}
$$

- Bounds can be obtained for $\left\|D g(y)^{-1}\right\|_{\infty}$ and $\|\Delta b\|$.
- Main idea: instead of the DFT,
- bound the WFT, and
- the difference WFT - DFT.


## Bound for $\left\|D g(y)^{-1}\right\|_{\infty}$

We can write

$$
D g(y)=: M=\left(\begin{array}{cccc}
2 & B_{0,1} & \ldots & B_{0, N_{f}} \\
0 & B_{1,1} & \ldots & B_{1, N_{f}} \\
\vdots & \vdots & \ddots & \vdots \\
0 & B_{N_{f}, 1} & \ldots & B_{N_{f}, N_{f}}
\end{array}\right)
$$

We split $M=M_{D}+M_{O}$,

$$
M=\left(\begin{array}{cccc}
2 & 0 & \ldots & 0 \\
0 & B_{1,1} & \ldots & 0 \\
\vdots & \vdots & \ddots & \vdots \\
0 & 0 & \ldots & B_{N_{f}, N_{f}}
\end{array}\right)+\left(\begin{array}{cccc}
0 & B_{0,1} & \ldots & B_{0, N_{f}} \\
0 & 0 & \ldots & B_{1, N_{f}} \\
0 & \vdots & \ddots & \vdots \\
0 & B_{N_{f}, 1} & \ldots & 0
\end{array}\right)
$$

$M$ is close to block-diagonal, so the idea is to obtain bounds for $\left\|M_{D}^{-1}\right\|,\left\|M_{O}\right\|$ and use

$$
\left\|\left(M_{D}+M_{O}\right)^{-1}\right\| \leq \frac{\left\|M_{D}^{-1}\right\|}{1-\left\|M_{D}^{-1}\right\|\left\|M_{O}\right\|}
$$

## Bound for $\|\Delta b\|_{\infty}$

We have

$$
\|\Delta b\| \leq 2 C \max _{j \in J} \sum_{|\boldsymbol{k}|=r_{0}}^{\infty} e^{-\delta|\boldsymbol{k}|}\left|\widetilde{h}_{N}^{n_{h}}(T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle-j)\right|
$$

where $\left|\widetilde{h}_{N}^{n_{h}}\right|$ is the envelope displayed below $\left(N=16, n_{h}=0\right)$.


## Bound for $\|\Delta b\|_{\infty}$

We have

$$
\|\Delta b\| \leq 2 C \max _{j \in J} \sum_{|\boldsymbol{k}|=r_{0}}^{\infty} e^{-\delta|\boldsymbol{k}|}\left|\widetilde{h}_{N}^{n_{h}}(T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle-j)\right|
$$

The Diophantine condition gives a lower bound for $|T\langle\boldsymbol{k}, \omega\rangle-j|$ :

$$
|T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle-j| \geq \frac{T D}{\left(|\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle|+\left|\boldsymbol{k}_{j}\right|\right)^{\tau}}-1 .
$$

For $|\boldsymbol{k}|$ small, $\left|\widetilde{h}_{N}^{n_{n}}(T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle-j)\right| \ll 1$.
After some order $r_{*},\left|\widetilde{h}_{N}^{n_{h}}(T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle-j)\right|$ may approach 1. Therefore,

$$
\|\Delta b\| \leq 2 C\left(\max _{j \in J} \sum_{|k|=r_{0}}^{r_{*}-1} e^{-\delta|k|}\left|\widetilde{h}_{N}^{n_{h}}(T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle-j)\right|+\max _{j \in J} \sum_{|k|=r_{*}}^{\infty} e^{-\delta|k|}\right) .
$$

## Bound for $\|\Delta b\|_{\infty}$

In

$$
\|\Delta b\| \leq 2 C\left(\max _{j \in J} \sum_{|\boldsymbol{k}|=r_{0}}^{r_{*}-1} e^{-\delta|\boldsymbol{k}|}\left|\widetilde{h}_{N}^{n_{h}}(T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle-j)\right|+\max _{j \in J} \sum_{|\boldsymbol{k}|=r_{*}}^{\infty} e^{-\delta|\boldsymbol{k}|}\right),
$$

- The first term is bounded by replacing the DFT by the WFT. This introduces an additional error term due to this approximation.
- All the sums are reduced to sums of the form $\sum_{j} j^{\alpha} e^{-\delta j}$, which are bounded by incomplete Gamma functions.


## Explicit bounds

Hypotheses:

1. Assume $f(t)=\sum_{\boldsymbol{k} \in \mathbb{Z}^{m}} a_{k} e^{i 2 \pi\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle t}$,

Cauchy estimates: $\left|a_{k}\right| \leq C e^{-\delta|k|}$,
$\boldsymbol{\omega}=\left(\omega_{1}, \ldots, \omega_{m}\right)$ rac ind.,
Diophantine condition $|\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle|>D /|k|^{\tau}$.
2. Apply the numerical Fourier analysis procedure with $T, N, n_{h}$ with minimum "amplitude barrier" $b_{\text {min }}$.
$\longrightarrow$ approximations $\widetilde{A}_{0},\left\{\left(\widetilde{\nu}_{k}, \widetilde{A}_{k}^{c}, \widetilde{A}_{k}^{s}\right)\right\}_{k=1}^{N_{f}}$
(denote by $A_{0},\left\{\left(\nu_{k}, A_{k}^{c}, A_{k}^{s}\right)\right\}_{k=1}^{N_{f}}$ the exact values)
3. Assume $\{T\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle\}_{|k|=1}^{r_{0}} \subset\left\{\nu_{k}\right\}_{k=1}^{N_{f}}$, for some order $r_{0}$,
4. $T, N$ satisfy some technical (non-demanding) lower bounds.

## Explicit bounds

Then the error can be bounded in first-order as:

$$
\|\Delta y\| \leq\left\|M^{-1}\right\|\|\Delta b\|
$$

with

$$
\begin{aligned}
& \text { - }\left\|M^{-1}\right\| \leq \frac{G_{n_{h}}}{\min \left(1, A_{\min }\right)}+\text { small terms } \quad \\
& \nabla\|\Delta b\| \leq \underbrace{\frac{C_{1}\left(n_{h}, m, C, \delta, D, \tau, r_{0}, r_{*}\right)}{T^{1+2 n_{h}}}}_{\text {leakage from orders } r_{0}, \ldots, r_{*}}+\underbrace{\frac{C_{2}\left(n_{h}, m, C, \delta, D, \tau, r_{0}, r_{*}\right)}{\left(D_{a}^{*}\right)^{1+2 n_{h}}}}_{\text {"aliasing" from orders } r_{0}, \ldots, r_{*}} \\
& +\underbrace{\left.\operatorname{tail}\left(n_{h}, m, C, \delta, D, \tau, r_{*}\right)\right)}_{\text {harmless amplitudes }}
\end{aligned}
$$

where $D_{a}^{*}:=N-T\left(r_{0}+r_{*}-2\right)\|\boldsymbol{\omega}\|_{\infty}-1$
is related to the distance of frequencies up to order $r_{*}$ to the right end of the fundamental domain of the DFT.

## Rules of Thumb for high accuracy

1. Choose $T$ such that the closest frequencies we want to determine are several harmonics away.
2. Choose $N$ such that the largest frequency we want to determine is away from the right end of the fundamental domain of the DFT.
3. Take $n_{h}=2$.

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## Outline

Introduction<br>The method<br>\section*{Error estimation}

## Accuracy test

Study of the stability region around $L_{5}$

## Accuracy test

We consider the quasi-periodic function $(\omega=(1, \sqrt{2}), \varphi=(0.2,0.3))$

$$
f_{\mu}(t)=\frac{\sin \left(2 \pi \omega_{1} t+\varphi_{1}\right)}{1-\mu \cos \left(2 \pi \omega_{1} t+\varphi_{1}\right)} \cdot \frac{\sin \left(2 \pi \omega_{2} t+\varphi_{2}\right)}{1-\mu \cos \left(2 \pi \omega_{2} t+\varphi_{2}\right)}, \quad \mu=0.9 .
$$

Explicit formulae for frequencies and amplitudes can be obtained, as well as the Cauchy estimates and the Diophantine condition.
We have performed Fourier analysis of this function for several $T, N$, computing the first 20 frequencies $(|k| \leq 5)$.



## Accuracy test

Error in amplitudes only:


For these functions, the Cauchy estimates are equalitites:

$$
f_{\mu}(t)=\sum_{\boldsymbol{k} \in \mathbb{Z}^{m}} a_{\boldsymbol{k}} e^{i 2 \pi\langle\boldsymbol{k}, \boldsymbol{\omega}\rangle t}, \quad m=2, \quad\left|a_{\boldsymbol{k}}\right|=\frac{1}{\mu^{2}} c^{|\boldsymbol{k}|}=1.23 \cdot(0.627)^{|\boldsymbol{k}|}
$$

For $|\boldsymbol{k}|=6,\left|a_{\boldsymbol{k}}\right|=6.06 \times 10^{-2}$, but we get nearly full double-precision accuracy in frequencies and amplitudes.

## Outline

## Introduction <br> The method

## Error estimation

## Accuracy test

Study of the stability region around $L_{5}$



## The circular, planar RTBP



Equation of motion:

$$
\begin{aligned}
\ddot{x}-2 \dot{y} & =\partial_{x} \Omega(x, y), \\
\ddot{y}+2 \dot{x} & =\partial_{y} \Omega(x, y),
\end{aligned}
$$

where

$$
\begin{aligned}
& r_{1}=\sqrt{(x-\mu)^{2}+y^{2}} \\
& r_{2}=\sqrt{(x-\mu+1)^{2}+y^{2}}
\end{aligned}
$$

$$
\Omega(x, y)=\frac{1}{2}\left(x^{2}+y^{2}\right)+\frac{1-\mu}{r_{1}}+\frac{\mu}{r_{2}}+\frac{1}{2} \mu(1-\mu) .
$$

Mass parameter: $\mu=\frac{m_{1}}{m_{1}+m_{2}}$.

## Data for the Sun-Jupiter case

- Sun-Jupiter mass parameter:

$$
\mu_{\mathrm{SJ}}=1 / 1048.3486=9.5388118 \times 10^{-4}
$$

- $L_{5}$ is center $\times$ center: $\quad \operatorname{Spec} D \boldsymbol{f}\left(L_{5}\right)=\left\{\omega_{\text {long }}^{L_{5}}, \omega_{\text {short }}^{L_{5}}\right\}$,

$$
\begin{aligned}
& \omega_{\text {long }}^{L_{5}}=\left(\frac{1-\sqrt{1-27 \mu(1-\mu)}}{2}\right)^{1 / 2}=0.08046412 \\
& \omega_{\text {short }}^{L_{5}}=\left(\frac{1+\sqrt{1-27 \mu(1-\mu)}}{2}\right)^{1 / 2}=0.99675750
\end{aligned}
$$

## Data for the Sun-Jupiter case

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$$
\omega_{\text {long }}^{L_{5}}=0.08046412, \quad \omega_{\text {short }}^{L_{5}}=0.99675750
$$

- We'll work with frequencies in cycles per unit of synodic time:

$$
\begin{aligned}
\nu_{\text {short }}^{L_{5}} & =\omega_{\text {short }}^{L_{5}} /(2 \pi) \\
\nu_{\text {long }}^{L_{5}} & =\omega_{\text {long }}^{L_{5}} /(2 \pi)
\end{aligned}=0.01280626, ~=0.15863888,
$$

- NOTE: $\nu_{\text {short }}^{L_{5}} / \nu_{\text {long }}^{L_{5}}=12.3876$.


## The stability domain

Numerical computation (G. Gómez, À. Jorba, J.J. Masdemont, C. Simó, ESA report 1993)


Parametrize the neighborhood of $L_{5}$ by

$$
\binom{x}{y}=\binom{\mu}{0}+(1+\rho)\binom{\cos (2 \pi \alpha)}{\sin (2 \pi \alpha)}
$$

For a grid of values of $\alpha, \rho$, take i.c.

$$
\begin{aligned}
& x_{0}=\mu+(1+\rho) \cos (2 \pi \alpha), \\
& y_{0}=(1+\rho) \sin (2 \pi \alpha), \\
& \dot{x}_{0}=\dot{y}_{0}=0 .
\end{aligned}
$$

Try to integrate up to time $T_{\text {max }}$, satisfying:

- Projection on $(x, y)$ not encircling the main primary.
- Not too close aproaches to primaries.
- $y>y_{c}=-0.5$.


## The stability domain

## Refinement (C. Simó, 2006, 2008)



- First run: up to $T_{\max }=2^{20}(2 \pi)$. Subsisting points: 215673.
- Second run: try the previous points up to $T_{\max }=2^{24}(2 \pi)$. Not all points are tested, but:
- From the border to the inside.
- Stop testing when 5 consecutive points stay for $2^{24}$ Jupiter revolutions.
Subsisting points: 215115.

Note: This is not the phase portrait on an area-preserving map. The initial conditions correspond to different energy levels.
Goal: to relate the frontier of the domain of stability and the island structure to resonances.

## The stability domain



## Fourier exploration

- The Fourier analysis procedure has been applied to each of the subsisting points, with

$$
T=65536, N=262144, n_{h}=2, N_{\max }=100, b_{\min }=10^{-6}
$$

- Total computing time: 352.52 hours (using 28 processors: 12.59 hours)
- Statistics:

|  | \#analyses |  |
| :--- | ---: | ---: |
| OK | 205779 | $95.41 \%$ |
| frequencies too close | 8722 | $4.04 \%$ |
| refinement did not converge | 878 | $0.41 \%$ |
| the two of the above | 294 | $0.14 \%$ |
| TOTAL | 215673 | $100 \%$ |

## Basic frequencies




- Left:
- Blue: freq. of maximum amplitude. It is close to $\nu_{\text {long }}^{L_{5}}$
$\longrightarrow \nu_{\text {long }}$
- Red: frequency of maximum amplitude inside [0.155, 0.165].

It is close to $\nu_{\text {short }}^{L_{5}}$
$\longrightarrow \nu_{\text {short }}$

- Right: the quotient $\nu_{\text {short }} / \nu_{\text {long }}$ for $\rho=4950$.


## Results

A basic set has been extracted from each set of frequencies, and all frequencies have been written as linear combinations of the basic set.
This allows to classify all the points in 4 groups:

1. Analyses ending with an error code. 9894 (4.54\%)
2. Error in determination of linear combinations $\geq 10^{-10}$. 20416 (9.47\%)
3. $\nu_{\text {short }}$ is not a rational multiple of $\nu_{\text {long }}$.

170389 (79.09\%)
4. $\nu_{\text {short }}$ is a rational multiple of $\nu_{\text {long }}$.

14914 (6.91\%)
$1+2$ : diffusing (chaotic) orbits.
3 : regular, non-resonant motion.
4 : regular, resonant motion.

## Graphical representation



- Blue: all the analyses
> Dark gray: ended with error code - Green:
error $>10^{-10}$ in determination of linear combinations
$\rightarrow$ Red:
$\nu_{\text {short }}$ not resonant with $\nu_{\text {long }}$


## Graphical representation



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Resonances: 14:1

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Resonances: 14:1, 29:2

## Graphical representation



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- Dark gray: ended with error code
- Green:
error $>10^{-10}$ in determination of linear combinations
- Red:
$\nu_{\text {short }}$ not resonant with $\nu_{\text {long }}$

Resonances: 14:1, 29:2, 15:1

## Graphical representation



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- Dark gray: ended with error code
- Green:
error $>10^{-10}$ in determination of linear combinations
- Red:
$\nu_{\text {short }}$ not resonant with $\nu_{\text {long }}$

Resonances: 14:1, 29:2, 15:1, 31:2

## Graphical representation



- Blue: all the analyses
- Dark gray: ended with error code
- Green:
error $>10^{-10}$ in determination of linear combinations
- Red:
$\nu_{\text {short }}$ not resonant with $\nu_{\text {long }}$

Resonances: 14:1, 29:2, 15:1, 31:2, 16:1

## Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2

## Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1

## Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2

## Graphical representation



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- Green:
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- Red:
$\nu_{\text {short }}$ not resonant with $\nu_{\text {long }}$

Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1

## Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2

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## \& that's it

Thank you!!

