Numerical Fourier analysis of quasi-periodic functions

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Outline

Introduction

The method

Error estimation

Accuracy test

Study of the stability region around L_5

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Setting

We are given an analytic, quasi-periodic function

$$f(t) = \sum_{\boldsymbol{k} \in \mathbb{Z}^m} a_{\boldsymbol{k}} e^{i2\pi \langle \boldsymbol{k}, \boldsymbol{\omega} \rangle t},$$

satisfying the Cauchy estimates

 $|a_{k}| \leq Ce^{-\delta|k|} \quad (\exists C > 0, \delta > 0, \quad |k| = |(k_{1}, \dots, k_{m})| = |k_{1}| + \dots + |k_{m}|)$

and with a vector of basic frequencies $\boldsymbol{\omega} = (\omega_1, \dots, \omega_m)$ satisfying a Diophantine condition

$$|\langle \boldsymbol{k}, \boldsymbol{\omega}
angle| > rac{D}{|k|^{ au}}, \quad igl(\exists D, au > 0 igr).$$

We want to **numerically compute** the frequencies $\{\langle k, \omega \rangle\}_{|k|=0}^{\max or}$ and amplitudes a_k from the values of f.

Fourier Transform

The Fourier Transform will be denoted as

$$f(t) \xrightarrow{\mathcal{F}} \mathcal{F}(f(t))(\omega) = \hat{f}(\omega) = \int_{-\infty}^{\infty} f(t)e^{-i2\pi\omega t}dt$$

If f(t) is quasi-periodic, its Fourier transform is a discrete set of impulses based at the frequencies:



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Example: $f(t) = \cos(2\pi 0.1t) + 0.5\cos(2\pi 0.2t) + 0.4\cos(2\pi 0.35t)$

Time truncation \longrightarrow WFT

Graphical development (E.O. Brigham, 1988)

Time truncation gives rise to the phenomenon known as *leakage*. Example: $T = 40, f(t) = \cos(2\pi 0.1t) + 0.5\cos(2\pi 0.2t) + 0.4\cos(2\pi 0.35t)$. 1.5 1.5 1.5 0.5 0.5 0.5 Х 0 0 0 -0.5 -0.5 -0.5 -1 -1 -1 -1.5 -1.5 -1.50 10 20 30 40 50 60 -20 -10 0 20 30 40 50 60 -20 -10 0 10 20 30 40 50 60 -20 -10 10 $\downarrow \mathcal{F}$ $|\mathcal{F}|$ $|\mathcal{F}|$ 1.2 40 35 35 30 30 0.8 25 25 0.6 20 20 0.4 15 10 0.2 5 0 MMM -0.2 -0.6 -0.4 -0.2 0.2 0.4 0.6 -0.6 -0.4 -0.2 0.4 0 0 0.2 0.6 -0.6 -0.4 -0.2 0 0.2 0.4 0.6

The maxima of the WFT (bottom right) are displaced from the true frequencies.

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Time truncation \longrightarrow WFT

Explicit formulae

Windowed Fourier Transform:

$$egin{aligned} \phi_{f,T}(\omega) &:= & rac{1}{T}\mathcal{F}ig(\chi_{[0,T]}f(t)ig)(\omega) \ &= & rac{1}{T}\int_0^T\chi_{[0,T]}(t)f(t)e^{-i2\pi\omega t}dt. \end{aligned}$$

• Leakage of a complex exponential term.

$$\begin{aligned} \phi_{e^{i2\pi\nu t},T}(\omega)| &= \left| \frac{e^{i2\pi(\nu-\omega)T} - 1}{i2\pi(\nu-\omega)T} \right| \\ &= \left| \frac{\sin \pi(\nu-\omega)T}{\pi(\nu-\omega)T} \right| \\ &= |\operatorname{sinc}((\nu-\omega)T)| \end{aligned}$$

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Reducing leackage

There are two strategies:

Increase the window length.

$$|\phi_{e^{i2\pi\nu \nu},T}(\omega)| = |\operatorname{sinc}((\nu-\omega)T)| = \left|\frac{\sin\pi(\nu-\omega)T}{\pi(\nu-\omega)T}\right|$$



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Reducing leackage

There are two strategies:

Use a smoother window.
 We use Hanning's:

$$H_T^{n_h}(t) = q_{n_h} \left(1 - \cos \frac{2\pi t}{T}\right)^{n_h}.$$

being
$$q_{n_h} = n_h!/((2n_h - 1)!!).$$

The corresponding WFT is denoted by

$$\phi_{f,T}^{n_h}(\omega) := \mathcal{F}ig(H^{n_h}fig)(\omega) = rac{1}{T}\int_0^T H_T^{n_h}(t)f(t)e^{-i2\pi\omega t}dt,$$

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Reducing leackage

There are two strategies:

► Use a smoother window.

$$\begin{split} \phi_{e^{i2\pi\nu t},T}(\omega) &= \frac{e^{i2\pi(\nu-\omega)T}-1}{i2\pi(\nu-\omega)T} = O\bigg(\frac{1}{(\nu-\omega)T}\bigg), \\ \phi_{e^{i2\pi\nu t},T}^{n_h}(\omega) &= \frac{(-1)^{n_h}(n_h!)^2 \big(e^{i2\pi(\nu-\omega)T}-1\big)}{i2\pi\prod_{j=-n_h}^{n_h}((\nu-\omega)T+j)} = O\bigg(\frac{1}{((\nu-\omega)T)^{1+2n_h}}\bigg) \end{split}$$



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Discretization \longrightarrow DFT

Graphical development (E.O. Brigham, 1988)



Sampling
$$\longrightarrow$$
 DFT

Explicit formulae

► DFT of $\{f(j_{\overline{N}}^T)\}_{j=0}^{N-1}$ defined as $\{F_{f,T,N}(k)\}_{k=0}^{N-1}$, being $F_{f,T,N}(k) := \frac{1}{N} \mathcal{F}\left(\sum_{j\in\mathbb{Z}} \chi_{[0,T]}\left(j\frac{T}{N}\right) f\left(j\frac{T}{N}\right) \delta_{j_{\overline{N}}^T}\right) \left(\frac{k}{T}\right)$ $= \frac{1}{N} \sum_{i=0}^{N-1} f\left(j\frac{T}{N}\right) e^{-i2\pi k j/N}.$

• With Hanning's window:

$$F_{f,T,N}^{n_h}(k) = rac{1}{N} \sum_{j=0}^{N-1} H_T^{n_h}(jrac{T}{N}) f(jrac{T}{N}) e^{-i2\pi k j/N}.$$

Sampling $\longrightarrow DFT$

Explicit formulae

Relation with the WFT:

$$F_{f,T,N}(k) = \phi_{f,T,N}\left(\frac{k}{T}\right) + \underbrace{\sum_{l \in \mathbb{Z} \setminus \{0\}} \left(\phi_{f,T,N}\left(\frac{k+lN}{T}\right) + \phi_{f,T,N}\left(\frac{k-lN}{T}\right)\right)}_{\text{error term}}$$

• The **fundamental domain** of the DFT for real signals is [0, T/(2N)]. T/(2N) is Nyquist's critical frequency.



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$Sampling \longrightarrow DFT$

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- The **fundamental domain** of the DFT for real signals is [0, T/(2N)]. T/(2N) is Nyquist's critical frequency.
- The error term above can produce aliasing: if a frequency of the signal is outside the fundamental domain of the DFT, we will detect an alias of it.

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► Aliasing is avoided increasing *N*.

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- Parameters: *T* (time length), *N* (number of samples), n_h (Hanning index) b_{\min} minimum threshold, several tolerances.
 - 1. Set an starting threshold for collecting peaks of the modulus of the DFT of f(t).
 - 2. Find initial **approximations of the frequencies**, starting from the peaks of the DFT greater than the thresold.
 - 3. Find the **amplitudes** of the frequencies found in the previous step, by solving $DFT(Q_f) = DFT(f)$.
 - Simultaneously refine ALL the frequencies and amplitudes of the current quasi-periodic approximation of *f*, by solving DFT(Q_f) = DFT(*f*).
 - 5. Perform a DFT of the input signal minus the current quasi-periodic approximation obtained in step 4, decrease the thresold and go back to step 2.

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For
$$f(t) = \cos(2\pi 0.13t) - \frac{1}{2}\sin(2\pi 0.27t) + \sin(2\pi 0.37t)$$
,
 $T = N = 512, n_h = 0.$

1. Starting thresold: 0.8 modulus of the DFT of the input data:



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For
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,
 $T = N = 512, n_h = 0.$

2. Approximation of frequencies:

peak 67 \Rightarrow frequency 0.130859375

peak 189 \Rightarrow frequency 0.369140625

3. Computation of amplitudes from known frequencies:

Frequency	Cosine amplitude	Sine amplitude
0.369140625	0.702312716711	0.136800713691
0.130859375	0.137731069235	0.699288924190

modulus of the DFT of the residual



= 900

For
$$f(t) = \cos(2\pi 0.13t) - \frac{1}{2}\sin(2\pi 0.27t) + \sin(2\pi 0.37t)$$
,
 $T = N = 512, n_h = 0.$

4. Iterative refinement:

Frequency	Cosine amplitude	Sine amplitude
0.369995932915	0.005462459021	1.000450861577
0.129998625183	0.999908805689	-0.002241420351

5. modulus of the DFT of input signal minus step 4:



New threshold: 0.2

For
$$f(t) = \cos(2\pi 0.13t) - \frac{1}{2}\sin(2\pi 0.27t) + \sin(2\pi 0.37t)$$
,
 $T = N = 512, n_h = 0.$

5. modulus of the DFT of input signal minus step 4:



New threshold: 0.2

2. Approximation of frequencies:

peak 138 \Rightarrow frequency 0.26953125

3. Amplitudes from known frequencies:

Frequency	Cosine amplitude	Sine amplitude
0.369995932915	0.005462459021	1.000450861577
0.129998625183	0.999908805689	-0.002241420352
0.269531250000	-0.309714556917	-0.330986794067

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For
$$f(t) = \cos(2\pi 0.13t) - \frac{1}{2}\sin(2\pi 0.27t) + \sin(2\pi 0.37t)$$
,
 $T = N = 512, n_h = 0.$

4. Iterative refinement:

 Frequency
 Cosine amplitude
 Sine amplitude

 0.370000000000000
 0.000000000000
 1.00000000000022

 0.13000000000000
 0.999999999999999
 0.00000000000000

 0.270000000000000
 -0.0000000000028
 -0.499999999999999

modulus of the DFT of the residual:



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Computing amplitudes from known frequencies We ask $DFT(Q_f) = DFT(f)$, being

$$Q_f(t) = A_0^c + \sum_{l=1}^{N_f} \left(A_l^c \cos(2\pi \frac{\nu_l}{T} t) + A_l^s \sin(2\pi \frac{\nu_l}{T} t) \right).$$

Since we work with **real** signals, we use the sine and cosine transforms:

$$\begin{aligned} c_{f,T,N}^{n_h}(k) &= \frac{2}{N} \sum_{j=0}^{N-1} f(j_N^T) H_N^{n_h}(j) \cos\left(2\pi \frac{k}{N}j\right), \quad k = 0, ..., \frac{N}{2}, \\ s_{f,T,N}^{n_h}(k) &= \frac{2}{N} \sum_{j=0}^{N-1} f(j_N^T) H_N^{n_h}(j) \sin\left(2\pi \frac{k}{N}j\right), \quad k = 1, ..., \frac{N}{2} - 1. \end{aligned}$$

They are realted to the DFT in complex form by

$$F_{f,T,N}^{n_h}(k) = \frac{1}{2} \Big(c_{f,T,N}^{n_h}(k) - i s_{f,T,N}^{n_h}(k) \Big), \qquad k = 0, \dots, N/2.$$

Computing amplitudes from known frequencies

The system of equations to be solved is **linear** and $(1 + 2N_f) \times (1 + 2N_f)$:

$$\begin{aligned} A_{0}^{c}c_{1,T,N}^{n_{h}}(0) + \sum_{l=1}^{N_{f}} & \left(A_{l}^{c}\overline{c}_{\nu_{l},N}^{n_{h}}(0) + A_{l}^{s}\widetilde{c}_{\nu_{l},N}^{n_{h}}(0)\right) &= c_{f,T,N}^{n_{h}}(0) \\ & A_{0}^{c}c_{1,T,N}^{n_{h}}(j) + \sum_{l=1}^{N_{f}} & \left(A_{l}^{c}\overline{c}_{\nu_{l},N}^{n_{h}}(j) + A_{l}^{s}\widetilde{c}_{\nu_{l},N}^{n_{h}}(j)\right) &= c_{f,T,N}^{n_{h}}(j) \\ & \sum_{l=1}^{N_{f}} & \left(A_{l}^{c}\overline{s}_{\nu_{l},T}^{n_{h}}(j) + A_{l}^{c}\widetilde{s}_{\nu_{l},T}^{n_{h}}(j)\right) &= s_{f,T,N}^{n_{h}}(j) \end{aligned}$$

where $j = [\nu_l + 0.5], l = 1 \div N_f$ (collocation harmonics), and

Simultaneous improvement of frequencies and amplitudes

We solve by Newton's method the following $(1 + 3N_f) \times (1 + 3N_f)$ non–linear system:

$$\begin{split} A_{0}^{c}c_{1,T,N}^{n_{h}}(0) + \sum_{l=1}^{N_{f}} \left(A_{l}^{c}\overline{c}_{\nu_{l},N}^{n_{h}}(0) + A_{l}^{s}\widetilde{c}_{\nu_{l},N}^{n_{h}}(0) \right) &= c_{f,T,N}^{n_{h}}(0) \\ A_{0}^{c}c_{1,T,N}^{n_{h}}(j_{i}) + \sum_{l=1}^{N_{f}} \left(A_{l}^{c}\overline{c}_{\nu_{l},N}^{n_{h}}(j_{i}) + A_{l}^{s}\widetilde{c}_{\nu_{l},N}^{n_{h}}(j_{i}) \right) &= c_{f,T,N}^{n_{h}}(j_{i}) \\ \sum_{l=1}^{N_{f}} \left(A_{l}^{c}\overline{s}_{\nu_{l},N}^{n_{h}}(j_{i}) + A_{l}^{s}\widetilde{s}_{\nu_{l},N}^{n_{h}}(j_{i}) \right) &= s_{f,T,N}^{n_{h}}(j_{i}) \\ A_{0}^{c}cs_{1,T,N}^{n_{h}}(j_{i}^{+}) + \sum_{l=1}^{N_{f}} \left(A_{l}^{c}\overline{c}\overline{s}_{\nu_{l},N}^{n_{h}}(j_{i}^{+}) + A_{l}^{s}\widetilde{c}\widetilde{s}_{\nu_{l},N}^{n_{h}}(j_{i}^{+}) \right) &= cs_{f,T,N}^{n_{h}}(j_{i}^{+}) \end{split}$$

being $j_i = [\nu_i + 0.5], j_i^+ = [\nu_i] + 1 - (j_i^+ - [\nu_i]).$

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Strategy

Let us denote

• f_{r_0} : the truncation of f to the frequencies we want to determine:

$$f_{r_0}(t) = A_0^c + \sum_{\substack{|\boldsymbol{k}| \le r_0 - 1 \ \langle \boldsymbol{k}, \boldsymbol{\omega}
angle > 0}} (A_k^c \cos(2\pi \langle \boldsymbol{k}, \boldsymbol{\omega}
angle t) + A_k^s \sin(2\pi \langle \boldsymbol{k}, \boldsymbol{\omega}
angle t)).$$

- ► $y = (A_0, \nu_1, A_1^c, A_1^s, \dots, \nu_{N_f}, A_{N_f}^c, A_{N_f}^s)$: the **exact** frequencies and amplitudes.
- $y + \Delta y$: the **computed** frequencies and amplitudes.

The **system we solve** for iterative improvement of frequencies and amplitudes is

$$\underbrace{\operatorname{DFT}(\mathcal{Q}_f)}_{g(y+\Delta y)} = \underbrace{\operatorname{DFT}(f_{r_0})}_{b} + \underbrace{\operatorname{DFT}(f-f_{r_0})}_{\Delta b}$$

We would get the **exact** frequencies and amplitudes if $\Delta b = 0$.

Strategy

System for iterative improvement of frequencies and amplitudes:

$$\begin{aligned} A_{0}^{c} + \sum_{l=1}^{N_{f}} \left(A_{l}^{c} \overline{c}_{\nu_{l},N}^{n_{h}}(0) + A_{l}^{s} \overline{c}_{\nu_{l},N}^{n_{h}}(0) \right) &= c_{f_{r_{0}},T,N}^{n_{h}}(0) + c_{f-f_{r_{0}},T,N}^{n_{h}}(0) \\ A_{0}^{c} c_{1}^{n_{h}}(j_{i}) + \sum_{l=1}^{N_{f}} \left(A_{l}^{c} \overline{c}_{\nu_{l},N}^{n_{h}}(j_{i}) + A_{l}^{s} \overline{c}_{\nu_{l},N}^{n_{h}}(j_{i}) \right) &= c_{f_{r_{0}},T,N}^{n_{h}}(j_{i}) + c_{f-f_{r_{0}},T,N}^{n_{h}}(j_{i}) \\ \sum_{l=1}^{N_{f}} \left(A_{l}^{c} \overline{s}_{\nu_{l},N}^{n_{h}}(j_{i}) + A_{l}^{s} \overline{s}_{\nu_{l},N}^{n_{h}}(j_{i}) \right) &= s_{f_{r_{0}},T,N}^{n_{h}}(j_{i}) + s_{f-f_{r_{0}},T,N}^{n_{h}}(j_{i}) \\ A_{0}^{c} cs_{1}^{n_{h}}(j_{i}^{+}) + \sum_{l=1}^{N_{f}} \left(A_{l}^{c} \overline{cs}_{\nu_{l},N}^{n_{h}}(j_{i}^{+}) + A_{l}^{s} \widetilde{cs}_{\nu_{l},N}^{n_{h}}(j_{i}^{+}) \right) &= cs_{f_{r_{0}},T,N}^{n_{h}}(j_{i}^{+}) + cs_{f-f_{r_{0}},T,N}^{n_{h}}(j_{i}^{+}). \end{aligned}$$
where $f - f_{r_{0}} = \sum_{|\mathbf{k}| \ge r_{0}} a_{\mathbf{k}} e^{i2\pi \langle \mathbf{k}, \omega \rangle t}.$

• The error term Δb consists of DFT

• of periodic terms with frequencies not being computed,

• evaluated in harmonics corresponding to frequencies being computed. Therefore, the error term Δb can be considered leakage of the remainder, $f - f_{r_0}$.

Strategy

• The error term Δb can be considered **leakage of the remainder**

$$\text{DFT}(f - f_{r_0}) = \sum_{|\mathbf{k}| \ge r_0} a_{\mathbf{k}} \text{DFT}(e^{i2\pi \langle \boldsymbol{\omega}, \mathbf{k} \rangle t})$$

• The effect of the terms of the remainder on the error Δb is

- The DFT of terms corresponding to **low–order frequencies**, $\{\langle k, \omega \rangle\}_{|k| \ge r_0}$, evaluated at the harmonics $\{j_i, j_i^+\}$, will be **small** if the harmonics $T\langle k, \omega \rangle$ are far from $\{j_i, j_i^+\}$. This can be achieved by increasing *T* as long as there is no aliasing.
- The DFT of terms corresponding to high–order frequencies may not be small (*T*(*k*, ω) can be made arbitrarily close to a *j_i* for large enough |*k*|). However, the corresponding amplitudes will be small due to the Cauchy estimates

$$|a_k| \leq C e^{-\delta |k|} \quad \forall k \in \mathbb{Z}^m,$$

so they will be harmless.

Bounding

The system we solve for iterative improvement of frequencies and amplitudes is

$$\underbrace{\mathrm{DFT}(Q_f)}_{g(y+\Delta y)} = \underbrace{\mathrm{DFT}(f_{r_0})}_{b} + \underbrace{\mathrm{DFT}(f-f_{r_0})}_{\Delta b}$$

We would get the **exact** frequencies and amplitudes if $\Delta b = 0$.

> The error in frequencies and amplitudes is given, at first order, by

$$\|\Delta y\|_{\infty} \le \|Dg(y)^{-1}\|_{\infty} \|\Delta b\|_{\infty}.$$

- **•** Bounds can be obtained for $||Dg(y)^{-1}||_{\infty}$ and $||\Delta b||$.
- Main idea: instead of the DFT,
 - bound the WFT, and
 - ▶ the difference WFT DFT.

Bound for $||Dg(y)^{-1}||_{\infty}$

We can write

$$Dg(y) =: M = \begin{pmatrix} 2 & B_{0,1} & \dots & B_{0,N_f} \\ 0 & B_{1,1} & \dots & B_{1,N_f} \\ \vdots & \vdots & \ddots & \vdots \\ 0 & B_{N_f,1} & \dots & B_{N_f,N_f} \end{pmatrix}$$

We split $M = M_D + M_O$,

$$M = \begin{pmatrix} 2 & 0 & \dots & 0 \\ 0 & B_{1,1} & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & B_{N_f,N_f} \end{pmatrix} + \begin{pmatrix} 0 & B_{0,1} & \dots & B_{0,N_f} \\ 0 & 0 & \dots & B_{1,N_f} \\ 0 & \vdots & \ddots & \vdots \\ 0 & B_{N_f,1} & \dots & 0 \end{pmatrix}$$

.

M is **close to block-diagonal**, so the idea is to obtain **bounds for** $||M_D^{-1}||$, $||M_O||$ and use

$$\|(M_D + M_O)^{-1}\| \le \frac{\|M_D^{-1}\|}{1 - \|M_D^{-1}\|\|M_O\|}$$

Bound for $\|\Delta b\|_{\infty}$

We have

$$\|\Delta b\| \leq 2C \max_{j\in J} \sum_{|m{k}|=r_0}^{\infty} e^{-\delta|m{k}|} |\widetilde{h}_N^{n_h}(T\langlem{k},m{\omega}
angle-j)|$$

where $|\tilde{h}_N^{n_h}|$ is the envelope displayed below ($N = 16, n_h = 0$).



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Bound for $\|\Delta b\|_{\infty}$

We have

$$\|\Delta b\| \leq 2C \max_{j \in J} \sum_{|\boldsymbol{k}|=r_0}^{\infty} e^{-\delta|\boldsymbol{k}|} |\widetilde{h}_N^{n_h}(T\langle \boldsymbol{k}, \boldsymbol{\omega} \rangle - j)|$$

The Diophantine condition gives a lower bound for $|T\langle k, \omega \rangle - j|$:

$$|T\langle m{k},m{\omega}
angle - j| \ \geq \ rac{TD}{(|\langle m{k},m{\omega}
angle| + |m{k}_j|)^ au} - 1.$$

For $|\mathbf{k}|$ small, $|\widetilde{h}_N^{n_h}(T\langle \mathbf{k}, \omega \rangle - j)| \ll 1$. After some order r_* , $|\widetilde{h}_N^{n_h}(T\langle \mathbf{k}, \omega \rangle - j)|$ may approach 1. Therefore,

$$\|\Delta b\| \leq 2C \Big(\max_{j\in J} \sum_{|k|=r_0}^{r_*-1} e^{-\delta|k|} |\widetilde{h}_N^{n_h}(T\langle k, oldsymbol{\omega}
angle -j)| + \max_{j\in J} \sum_{|k|=r_*}^{\infty} e^{-\delta|k|} \Big).$$

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Bound for $\|\Delta b\|_{\infty}$

In

$$\|\Delta b\| \leq 2C \Big(\max_{j \in J} \sum_{|\boldsymbol{k}|=r_0}^{r_*-1} e^{-\delta|\boldsymbol{k}|} |\widetilde{h}_N^{n_h}(T\langle \boldsymbol{k}, \boldsymbol{\omega}\rangle - j)| + \max_{j \in J} \sum_{|\boldsymbol{k}|=r_*}^{\infty} e^{-\delta|\boldsymbol{k}|} \Big),$$

- The first term is bounded by replacing the DFT by the WFT. This introduces an additional error term due to this approximation.
- ► All the sums are reduced to sums of the form $\sum_j j^{\alpha} e^{-\delta j}$, which are bounded by incomplete Gamma functions.

Explicit bounds

Hypotheses:

- 1. Assume $f(t) = \sum_{k \in \mathbb{Z}^m} a_k e^{i2\pi \langle k, \omega \rangle t}$, Cauchy estimates: $|a_k| \leq Ce^{-\delta |k|}$, $\omega = (\omega_1, \dots, \omega_m)$ rac ind., Diophantine condition $|\langle k, \omega \rangle| > D/|k|^{\tau}$.
- 2. Apply the numerical Fourier analysis procedure with *T*, *N*, *n_h* with minimum "amplitude barrier" *b*_{min}.
 → approximations *A*₀, {(*v*_k, *A*^c_k, *A*^s_k)}^{N_f}_{k=1} (denote by *A*₀, {(*v*_k, *A*^c_k, *A*^s_k)}^k_{k=1} the exact values)
- 3. Assume $\{T\langle \boldsymbol{k}, \boldsymbol{\omega} \rangle\}_{|\boldsymbol{k}|=1}^{r_0} \subset \{\nu_k\}_{k=1}^{N_f}$, for some order r_0 ,
- 4. T, N satisfy some technical (non-demanding) lower bounds.

Explicit bounds

Then the error can be bounded in first-order as:

$$\|\Delta y\| \le \|M^{-1}\| \|\Delta b\|,$$

with

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Rules of Thumb for high accuracy

- 1. Choose *T* such that the closest frequencies we want to determine are several harmonics away.
- 2. Choose *N* such that the largest frequency we want to determine is away from the right end of the fundamental domain of the DFT.

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3. Take $n_h = 2$.

Rules of Thumb for high accuracy

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Rules of Thumb for high accuracy

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Outline

Introduction

The method

Error estimation

Accuracy test

Study of the stability region around L_5

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Accuracy test

We consider the **quasi-periodic function** ($\omega = (1, \sqrt{2}), \varphi = (0.2, 0.3)$)

$$f_{\mu}(t) = \frac{\sin(2\pi\omega_{1}t + \varphi_{1})}{1 - \mu\cos(2\pi\omega_{1}t + \varphi_{1})} \cdot \frac{\sin(2\pi\omega_{2}t + \varphi_{2})}{1 - \mu\cos(2\pi\omega_{2}t + \varphi_{2})}, \quad \mu = 0.9.$$

Explicit formulae for frequencies and amplitudes can be obtained, as well as the **Cauchy estimates** and the **Diophantine condition**.

We have performed **Fourier analysis** of this function for several *T*, *N*, computing the first 20 frequencies ($|k| \le 5$).



Accuracy test

Error in amplitudes only:



For these functions, the Cauchy estimates are equalitites:

$$f_{\mu}(t) = \sum_{k \in \mathbb{Z}^m} a_k e^{i2\pi \langle k, \omega \rangle t}, \quad m = 2, \quad |a_k| = \frac{1}{\mu^2} c^{|k|} = 1.23 \cdot (0.627)^{|k|}$$

For $|\mathbf{k}| = 6$, $|a_{\mathbf{k}}| = 6.06 \times 10^{-2}$, but we get nearly full double-precision accuracy in frequencies and amplitudes.

Outline

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The circular, planar RTBP



Equation of motion:

$$\begin{split} \ddot{x} - 2\dot{y} &= \partial_x \Omega(x,y), \\ \ddot{y} + 2\dot{x} &= \partial_y \Omega(x,y), \end{split}$$

where

$$r_1 = \sqrt{(x-\mu)^2 + y^2},$$

$$r_2 = \sqrt{(x-\mu+1)^2 + y^2},$$

$$\Omega(x,y) = \frac{1}{2}(x^2 + y^2) + \frac{1-\mu}{r_1} + \frac{\mu}{r_2} + \frac{1}{2}\mu(1-\mu).$$

Mass parameter: $\mu = \frac{m_1}{m_1 + m_2}$.

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Data for the Sun–Jupiter case

Sun–Jupiter mass parameter:

$$\mu_{\rm SJ} = 1/1048.3486 = 9.5388118 \times 10^{-4}$$

► L_5 is center × center: Spec $Df(L_5) = \{\omega_{\text{log}}^{L_5}, \omega_{\text{short}}^{L_5}\},\$

$$\omega_{\text{long}}^{L_5} = \left(\frac{1 - \sqrt{1 - 27\mu(1 - \mu)}}{2}\right)^{1/2} = 0.08046412,$$

$$\omega_{\text{short}}^{L_5} = \left(\frac{1 + \sqrt{1 - 27\mu(1 - \mu)}}{2}\right)^{1/2} = 0.99675750.$$

Data for the Sun–Jupiter case

Sun–Jupiter mass parameter:

$$\mu_{\rm SJ} = 1/1048.3486 = 9.5388118 \times 10^{-4}$$

► L_5 is center × center: Spec $Df(L_5) = \{\omega_{\text{long}}^{L_5}, \omega_{\text{short}}^{L_5}\},\$

$$\omega_{\text{long}}^{L_5} = 0.08046412, \qquad \omega_{\text{short}}^{L_5} = 0.99675750.$$

▶ We'll work with frequencies in cycles per unit of synodic time:

$$\begin{array}{rcl} \nu_{\rm short}^{L_{\rm S}} &=& \omega_{\rm short}^{L_{\rm S}}/(2\pi) &=& 0.01280626, \\ \nu_{\rm long}^{L_{\rm S}} &=& \omega_{\rm long}^{L_{\rm S}}/(2\pi) &=& 0.15863888, \end{array}$$

• NOTE: $\nu_{\text{short}}^{L_5} / \nu_{\text{long}}^{L_5} = 12.3876.$

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The stability domain

Numerical computation (G. Gómez, À. Jorba, J.J. Masdemont, C. Simó, ESA report 1993)



Parametrize the neighborhood of L_5 by

$$\left(\begin{array}{c} x\\ y\end{array}\right) = \left(\begin{array}{c} \mu\\ 0\end{array}\right) + (1+\rho) \left(\begin{array}{c} \cos(2\pi\alpha)\\ \sin(2\pi\alpha)\end{array}\right)$$

For a grid of values of α , ρ , take i.c.

$$\begin{aligned} x_0 &= \mu + (1+\rho)\cos(2\pi\alpha), \\ y_0 &= (1+\rho)\sin(2\pi\alpha), \\ \dot{x}_0 &= \dot{y}_0 = 0 \end{aligned}$$

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Try to integrate up to time T_{max} , satisfying:

- ▶ Projection on (*x*, *y*) not encircling the main primary.
- Not too close aproaches to primaries.

►
$$y > y_c = -0.5$$
.

The stability domain

Refinement (C. Simó, 2006, 2008)



- First run: up to $T_{\text{max}} = 2^{20}(2\pi)$. Subsisting points: 215673.
- Second run: try the previous points up to $T_{\text{max}} = 2^{24}(2\pi)$. Not all points are tested, but:
 - From the border to the inside.
 - Stop testing when 5 consecutive points stay for 2²⁴ Jupiter revolutions.

Subsisting points: 215115.

Note: This is **not** the phase portrait on an area-preserving map. The initial conditions correspond to different energy levels.

Goal: to relate the frontier of the domain of stability and the island structure to resonances.

The stability domain



Fourier exploration

The Fourier analysis procedure has been applied to each of the subsisting points, with

$$T = 65536, N = 262144, n_h = 2, N_{max} = 100, b_{min} = 10^{-6}$$

- Total computing time: 352.52 hours (using 28 processors: 12.59 hours)
- Statistics:

status	#analyses	
ОК	205 779	95.41%
frequencies too close	8 7 2 2	4.04%
refinement did not converge	878	0.41%
the two of the above	294	0.14%
TOTAL	215 673	100%

Basic frequencies



► Left:

• Blue: freq. of maximum amplitude. It is close to $\nu_{long}^{L_5}$

 $\longrightarrow \nu_{\text{long}}$

 Red: frequency of maximum amplitude inside [0.155, 0.165]. It is close to ν^{L₅}_{short}

- $\longrightarrow \nu_{\text{short}}$
- Right: the quotient $\nu_{\text{short}}/\nu_{\text{long}}$ for $\rho = 4950$.

Results

A basic set has been extracted from each set of frequencies, and all frequencies have been written as linear combinations of the basic set. This allows to classify all the points in 4 groups:

- 1. Analyses ending with an error code. 9894 (4.54%)
- 2. Error in determination of linear combinations $\geq 10^{-10}$. 20416 (9.47%)
- 3. ν_{short} is not a rational multiple of ν_{long} . 170389 (79.09%)
- 4. ν_{short} is a rational multiple of ν_{long} . 14914 (6.91%)
- 1 + 2 : diffusing (chaotic) orbits.
- 3 : regular, non-resonant motion.
- 4 : regular, resonant motion.

Graphical representation



Graphical representation



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Graphical representation



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Graphical representation



Graphical representation



Resonances: 14:1

Graphical representation



Resonances: 14:1, 29:2

Graphical representation



Resonances: 14:1, 29:2, 15:1



Resonances: 14:1, 29:2, 15:1, 31:2

Graphical representation



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Resonances: 14:1, 29:2, 15:1, 31:2, 16:1

Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2

Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2

Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2, 19:1

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Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2, 19:1, 39:2

Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2, 19:1, 39:2, 20:1

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Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2, 19:1, 39:2, 20:1, 41:2

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Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2, 19:1, 39:2, 20:1, 41:2, 21:1

Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2, 19:1, 39:2, 20:1, 41:2, 21:1, 22:1

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Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2, 19:1, 39:2, 20:1, 41:2, 21:1, 22:1, 23:1

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Graphical representation



Resonances: 14:1, 29:2, 15:1, 31:2, 16:1, 33:2, 17:1, 35:2, 18:1, 37:2, 19:1, 39:2, 20:1, 41:2, 21:1, 22:1, 23:1, 24:1

& that's it

Thank you!!